

Provided for non-commercial research and education use.  
Not for reproduction, distribution or commercial use.



Volume 403

Issues 5–9

1 April 2008

ISSN 0921-4526

**PHYSICA**



Recognized by the European Physical Society

**B**

**CONDENSED MATTER**



Proceedings of the International  
Conference on Strongly Correlated  
Electron Systems

**SCES 2007**

held in Houston, Texas, USA  
13–18 May 2007

Guest Editors:

B. Lorenz  
I. Vekhter  
C.J. Bolech  
P.C.W. Chu  
Q. Si

Available online at

 **ScienceDirect**  
www.sciencedirect.com

<http://www.elsevier.com/locate/physb>

This article was published in an Elsevier journal. The attached copy is furnished to the author for non-commercial research and education use, including for instruction at the author's institution, sharing with colleagues and providing to institution administration.

Other uses, including reproduction and distribution, or selling or licensing copies, or posting to personal, institutional or third party websites are prohibited.

In most cases authors are permitted to post their version of the article (e.g. in Word or Tex form) to their personal website or institutional repository. Authors requiring further information regarding Elsevier's archiving and manuscript policies are encouraged to visit:

<http://www.elsevier.com/copyright>



# Observation of two-magnon bound states in the spin-1 anisotropic Heisenberg antiferromagnetic chain system $\text{NiCl}_2-4\text{SC}(\text{NH}_2)_2$

S.A. Zvyagin<sup>a,\*</sup>, C.D. Batista<sup>b</sup>, J. Krzystek<sup>c</sup>, V.S. Zapf<sup>d</sup>,  
M. Jaime<sup>d</sup>, A. Paduan-Filho<sup>e</sup>, J. Wosnitza<sup>a</sup>

<sup>a</sup>Dresden High Magnetic Field Laboratory (HLD), Forschungszentrum Dresden—Rossendorf, 01314 Dresden, Germany

<sup>b</sup>Theoretical Division, Los Alamos National Laboratory, Los Alamos, NM 87545, USA

<sup>c</sup>National High Magnetic Field Laboratory, Florida State University, Tallahassee, FL 32310, USA

<sup>d</sup>National High Magnetic Field Laboratory, Los Alamos National Laboratory, MS-E536, Los Alamos, NM 87545, USA

<sup>e</sup>Instituto de Física, Universidade de Sao Paulo, 05315 970 Sao Paulo, Brazil

## Abstract

Results of systematic tunable-frequency ESR studies of the spin dynamics in  $\text{NiCl}_2-4\text{SC}(\text{NH}_2)_2$  (known as DTN), a gapped  $S = 1$  chain system with easy-plane anisotropy dominating over the exchange coupling (large- $D$  chain), are presented. We have obtained direct evidence for two-magnon bound states, predicted for  $S = 1$  large- $D$  spin chains in the fully spin-polarized (FSP) phase. The frequency–field dependence of the corresponding excitations was calculated using the set of parameters obtained earlier [S.A. Zvyagin, et al., Phys. Rev. Lett. 98 (2007) 047205]. Very good agreement between the calculations and the experiment was obtained. It is argued that the observation of transitions from the ground to two-magnon bound states might indicate a more complex picture of magnetic interactions in DTN, involving a finite in-plane anisotropy.

© 2007 Elsevier B.V. All rights reserved.

PACS: 75.40.Gb; 76.30.-v; 75.10.Jm

Keywords: Two-magnon bound states; Electron spin resonance; Field-induced phase transition; Bose–Einstein condensation

Antiferromagnetic (AFM) quantum spin-1 chains have been the subject of intensive theoretical and experimental studies, fostered especially by the Haldane conjecture [1]. The presence of a strong easy-plane anisotropy  $D$  can significantly modify the excitation spectrum [2] so that the gap size is not determined by the strength of the AFM quantum fluctuations *exclusively*, but depends on the dimensionless parameter  $\rho = D/J$  (where  $J$  is the exchange coupling). For  $\rho > 0.93$  [3], the origin of the gap is dominated by the anisotropy  $D$  and the system is in the so-called large- $D$  regime. The gap can be closed by a magnetic field,  $H_{c1}$ , applied perpendicular to the anisotropy plane. When the magnetic field exceeds the upper critical field,  $H_{c2}$ , the system is in the fully spin-polarized (FSP) state and the low-energy magnetic excitation

spectrum is formed by magnons. In addition to ordinary single-magnon states, the theory [4] predicts the existence of two-magnon bound states. In this paper, we report a detailed frequency–field diagram of magnetic excitations in DTN, a system of weakly interacting  $S = 1$  chains with single-ion anisotropy  $D$  larger than the intra-chain exchange coupling  $J_c$ , obtained by means of the high-field electron spin resonance (ESR) technique [5].

The compound  $\text{NiCl}_2-4\text{SC}(\text{NH}_2)_2$  (dichloro-*tetrakis* thiourea-nickel(II), known as DTN) has a tetragonal crystal structure with space group I4 [6]. The  $g$ -factor, anisotropy, intra- and inter-chain exchange parameters,  $g_c = 2.26$ ,  $D = 8.9$  K,  $J_c = 2.2$  K, and  $J_{a,b} = 0.18$  K, respectively, were obtained from high-field ESR data [7], zero-field inelastic-neutron scattering, and magnetocaloric measurements [8]. At sufficiently low temperatures ( $T < 1.2$  K) and at  $H \parallel c$  (where  $c$  is the tetragonal axis), DTN undergoes a transition into a field-induced AFM

\*Corresponding author. Tel.: +49 351 260 3517; fax: +49 351 260 3531.  
E-mail address: s.zvyagin@fzd.de (S.A. Zvyagin).

ordered phase with  $H_{c1} = 2.16$  T and  $H_{c2} = 12.6$  T (defined at  $T \rightarrow 0$ ) and critical exponents consistent with the Bose–Einstein-condensation universality class [8,9].

At  $H < H_{c1}$ , DTN is in the quantum-paramagnetic phase, having an  $S_z = 0$  ground state and a gapped excitation spectrum determined by the  $\Delta S_z = \pm 1$  transitions (the modes A and B in Fig. 1). Excitations from the ground to single-magnon states (the mode C in Fig. 1) were observed at  $T = 1.6$  K in the FSP phase. These excitations correspond to a single spin flip from the  $S_z = -1$  to 0 state and are uniformly delocalized over the entire lattice with a well-defined momentum  $\mathbf{k}$ . The ESR transitions taking place at  $\mathbf{k} = 0$  have the frequency  $\omega_C = g\mu_B H - D$ . From this expression, the single-ion anisotropy in DTN can be calculated accurately ( $D = 8.9$  K) [7].

The two-magnon bound states (sometimes referred to as single-ion bound states) were predicted in 1970 by Silbergliitt and Torrance [10] for Heisenberg ferromagnets with single-ion anisotropy. Later on, this subject attracted a great deal of attention due to its potential relevance to the intrinsic localized spin modes in anisotropic ferromagnets [11] and antiferromagnets [12]. The two-magnon bound states for  $S = 1$  AFM chains in the FSP phase have been predicted by Papanicolaou et al. [4]. The physical picture of the two-magnon bound-state excitations corresponds to a double-spin-flip transition from  $S_z = -1$  to  $+1$ . Since the diagonal energy difference between these two states,  $2D$ , is much bigger than exchange interactions, the distance between the two  $S_z = 0$  sites remains finite, giving rise to a two-magnon bound state. A first signature of two-magnon bound states was obtained by means of high-field ESR in the spin-1 chain compound  $\text{Ni}(\text{C}_2\text{H}_8\text{N}_2)_2\text{Ni}(\text{CN})_4$  (known as NENC) [14], which was interpreted as transitions from the single-magnon to two-magnon bound states.

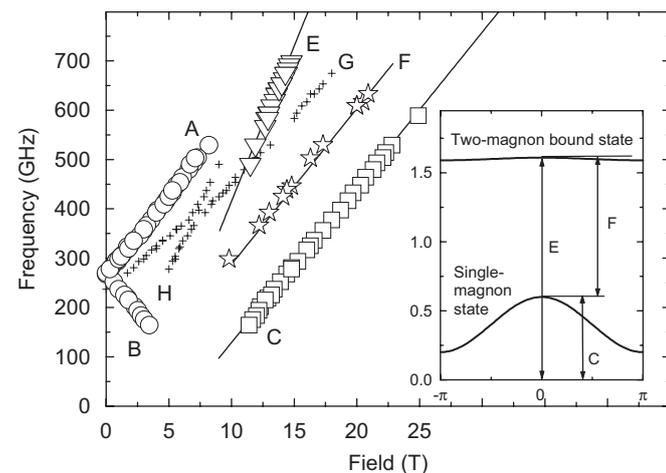


Fig. 1. Frequency–field dependence of magnetic excitations in DTN, taken at  $H \parallel c$ . Symbols denote the experimental results, and lines correspond to results of calculations (see text for details). The inset shows schematically the excitation dispersion in an  $S = 1$  Heisenberg chain with strong easy-plane anisotropy in the FSP phase (note that the ESR transitions denoted by C, E, and F occur at  $\mathbf{k} = 0$ ). The two-particle continuum is not shown for simplicity.

In this paper we report the direct observation of transitions from the ground state to the two-magnon bound state (the mode E, Fig. 1) in a spin-1 AFM chain system with strong easy-plane anisotropy in the high-field FSP phase. The corresponding excitations observed at  $T = 1.6$  K are denoted by triangles in Fig. 1 together with results of analytical calculations, using the set of parameters from Ref. [7] and the expression

$$\omega_E = -4(J_c + 2J_a) + 2J_c^2/D + 4J_a^2/D + 2g\mu_B H. \quad (1)$$

Very good agreement with experimental data was achieved. Since these excitations involve  $\Delta S_z = 2$  transitions, which are strictly forbidden in case of an ideal axial symmetry [13], their observation by means of ESR might indicate that the axial symmetry in FSP phase is broken.

Observation of the modes G and H (Fig. 1) might indicate existence of a second  $S = 1$  paramagnetic center in our samples. Although the exact origin of these modes is not clear at the moment, one can speculate that they are probably associated with excitations of  $\text{Ni}^{2+}$  spins in a superficial layer of DTN-crystals attacked by the GEvarnish solvent (used to fix the sample inside the sample-holder).

Finally, we observe a resonance absorption denoted by stars in Fig. 1. This ESR mode occurs at higher temperatures ( $T = 4.3$  K) and results from transitions within excited states, i.e. between the single-magnon and two-magnon bound states at  $\mathbf{k} = 0$ . The frequency–field dependence of this transition (the mode F) can be calculated using the expression  $\omega_F = \omega_E - \omega_C$  (where  $\omega_E$  and  $\omega_C$  are the excitation frequencies for the modes E and C, respectively).

In summary, an ESR study of the magnetic excitations in DTN, an  $S = 1$  Heisenberg AFM chain material with strong single-ion anisotropy, has been presented. Two-magnon bound states, predicted for  $S = 1$  large- $D$  spin chains in the FSP phase, have been observed directly. The observation of  $\Delta S_z = 2$  transitions from the ground to two-magnon bound states might indicate a more complex picture of the magnetic interactions in DTN, involving a finite single-ion in-plane anisotropy or the Dzyaloshinskii–Moriya interaction.

The authors express their sincere thanks to A.K. Kolezhuk and A. Orendáčová for fruitful discussions. S.A.Z. acknowledges the support from the NHMFL (which is supported by NSF Cooperative Agreement no. DMR-0084173, by the State of Florida, and by the DOE) through the VSP no. 1382.

## References

- [1] F.D.M. Haldane, Phys. Lett. 93A (1983) 464.
- [2] O. Golinelli, et al., Phys. Rev. B 46 (1992) 10854.
- [3] T. Sakai, M. Takahashi, Phys. Rev. B 42 (1990) 4537.
- [4] N. Papanicolaou, et al., Phys. Rev. B 56 (1997) 8786.
- [5] S.A. Zvyagin, et al., Physica B 346–347 (2004) 1.

- [6] A. Paduan-Filho, et al., *J. Chem. Phys.* 74 (1981) 4103.
- [7] S.A. Zvyagin, et al., *Phys. Rev. Lett.* 98 (2007) 047205.
- [8] V.S. Zapf, et al., *Phys. Rev. Lett.* 96 (2006) 077204.
- [9] A. Paduan-Filho, et al., *Phys. Rev. B* 69 (2004) 020405(R).
- [10] R. Silbergliitt, J.B. Torrance Jr., *Phys. Rev. B* 2 (1970) 772.
- [11] R.F. Wallis, et al., *Phys. Rev. B* 52 (1995) R3828.
- [12] R. Lai, et al., *Phys. Rev. B* 54 (1996) R12665.
- [13] A.K. Kolezhuk, et al., *Phys. Rev. B* 65 (2001) 014413.
- [14] S.A. Zvyagin, et al., *Czech. J. Phys.* 46 (1996) 1937;  
M. Orendáč, et al., *Phys. Rev. B* 60 (1999) 4170.