

# Fermi surface spectroscopy: a magnetic resonance approach

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## Abstract

We describe recently-developed GHz magneto-optical techniques which measure the Fermi-surface topologies and superconducting order parameter symmetry in organic molecular metals. © 2000 Elsevier Science B.V. All rights reserved.

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Quasi-two-dimensional (Q2D) organic molecular metals have attracted enormous attention in recent years, as they exhibit a wide range of low-temperature ground states and yet possess Q2D Fermi surfaces (FSs) which may be accurately represented using tractable analytical models [1]. Here we review recently derived methods for establishing the FS topology of these systems which involve GHz magneto-optical techniques.

The path of a quasiparticle of energy  $E(\mathbf{k})$  and velocity  $\mathbf{v}$  in a magnetic field  $\mathbf{B}$  is determined by the equations  $\hbar\mathbf{v} = \nabla_{\mathbf{k}}E(\mathbf{k})$  and  $\hbar\partial\mathbf{k}/\partial t = -e\mathbf{v} \times \mathbf{B}$ , where  $\nabla_{\mathbf{k}}$  is the gradient operator in  $k$ -space. Quasiparticles at the FS

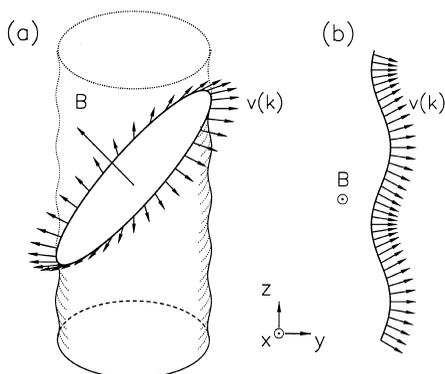


Fig. 1. Schematic view of two types of orbit on the Fermi surfaces (FSs) of organic metals in a magnetic field  $\mathbf{B}$ . (a) closed orbit around warped Q2D FS section. (b) open orbit along Q1D FS section. The arrows indicate the quasiparticle velocities.

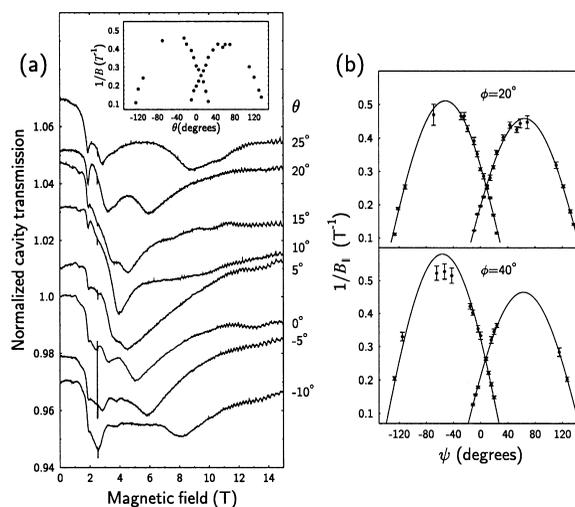


Fig. 2. (a) FTRs in  $\alpha$ -(ET)<sub>2</sub>KHg(SCN)<sub>4</sub> at  $T = 1.4$  K and 70 GHz using a rectangular cavity tilted to angles  $\theta$  in the field. FTR fields are plotted versus  $\theta$  (inset) and  $\psi$  (b), where  $\psi$  is the angle between  $\mathbf{B}$  and the warping direction of the Q1D FS. Data are shown in (b) for two orientations of the sample within the cavity. The  $\psi$  dependence in (b) unambiguously identifies FTRs [4].

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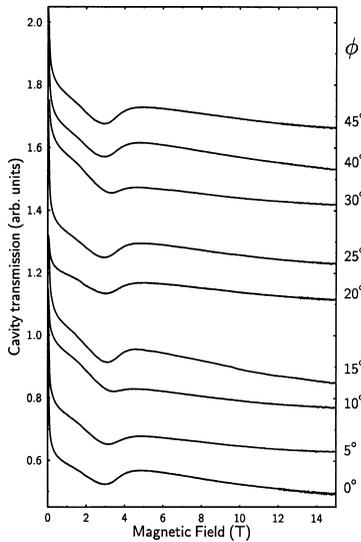


Fig. 3. Transmission of a cavity containing a  $\kappa$ -(ET)<sub>2</sub>Cu(NCS)<sub>2</sub> crystal (70 GHz, 1.4 K); the Q2D planes are parallel to  $\vec{H}$  and normal to  $\vec{B}$ , but  $\vec{H}$  is rotated in the Q2D plane by angles  $\phi$ ; the induced in-plane current thus varies in direction. The dip at  $\sim 3$  T marks the change in dissipation at  $H_{c2}$ ; its size varies with  $\phi$ , reflecting the angular variation of the order parameter [5].

will move on orbits of constant energy in the plane perpendicular to  $\vec{B}$ . Fig. 1 shows two classes of orbit important in organic molecular metals; (a) is a closed orbit about a Q2D FS section and (b) is an open orbit along a quasi-one-dimensional (Q1D) FS sheet. Note that the velocities of the quasiparticles (indicated by arrows in Fig. 1) oscillate up and down (a) or from side to side (b); the Boltzmann transport equation has been used to show that these oscillations contribute resonantly to the bulk ac conductivity [2,3]. As the corrugations of real FSs will not in general be sinusoidal, the conductivity will also contain a series of harmonics of the fundamental frequency which in effect represent a “Fourier analysis” of the shape of the Fermi-surface section [2,3]. Resonances from Q2D sections of FS are termed “periodic orbit resonances” (PORs) [2]; those from the Q1D FS sections are called Fermi-surface traversal resonances (FTRs) [4]. The high-frequency bulk conductivity is mea-

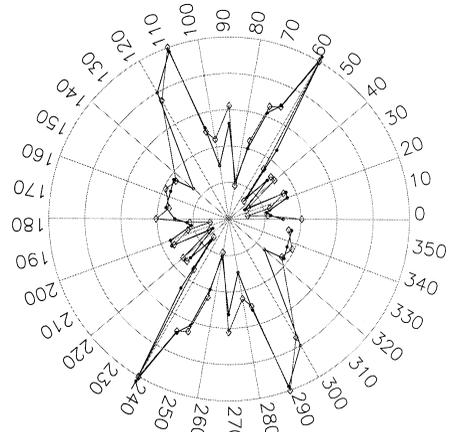


Fig. 4. The normalised amplitude of the dip in Fig. 3 versus  $\phi$ ; The underlying d-like symmetry of the order parameter can clearly be seen [1,5].

sured using the transmission of a rectangular [4] or cylindrical [2] resonant cavity loaded with a sample. The oscillating  $H$ -field  $\vec{H}$  in the cavity is parallel to the sample’s highly conducting Q2D planes; induced currents flow in the low-conductivity interplane direction, so that the skin-depth  $\gg$  the sample size [4]. FTRs have a characteristic dependence on the direction of  $\vec{B}$ ; hence, experiments are done using rectangular cavities which tilt in the static field [4]. FTRs observed in the organic metal  $\alpha$ -(ET)<sub>2</sub>KHg(SCN)<sub>4</sub> are seen in Fig. 2 as broad absorptions above  $\sim 4$  T [4].

The bulk GHz conductivity is also affected by the presence of superconductivity. GHz conductivity measurements have been used to identify the presence of an anisotropic (d-wave-like) order parameter in the organic superconductor  $\kappa$ -(ET)<sub>2</sub>Cu(NCS)<sub>2</sub> (see Figs. 3 and 4) [5].

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